

ordination has only been described for $\text{RuCl}_2\text{C}_{12}\text{H}_{18}$, where the organic ligand is the linear trimer of butadiene (Lydon, Nicholson, Shaw & Truter, 1964; Lydon & Truter, 1968). The purpose of this paper is to give the details of the X-ray structure of (I).

Experimental

Red-brown, well formed, air stable single crystals of (I) were obtained by maintaining a solution of RuCl_3 in ethanol (0.1g in 2 cm³) with a large excess (10 cm³) of isoprene at a temperature of 60–70°C for 40 hours (Porri *et al.*, 1965), and then slowly cooling the reacted mixture.

X-ray intensities were collected with the multiple-film equi-inclination Weissenberg technique, and Mo $K\alpha$ radiation. The integrated intensities of 1124 non-zero reflexions were estimated visually on the ($hk0$), ($h0l$), ($h1l$), ($0kl$) and ($1kl$) layers. No absolute scaling of the reflexions was performed at this stage; the scaling factors of the various layers were determined subsequently as adjustable parameters during the least-squares refinement. Eventually, the structure factors of the reflexions common to different layers were averaged; the resulting number of distinct reflexions was 1008. No absorption factor was applied, on account of the small size and of the regular, cubic shape of the crystal ($\mu R < 1$). The parameters of the triclinic unit cell were measured on the three equatorial layer Weissenberg photographs, calibrated with a powder spectrum of LiF; they are reported in Table 1.

Table 1. *Crystal data*

$(\text{RuCl}_2\text{C}_{10}\text{H}_{16})_2$ M.W. 616.5

| | |
|-------------------|---------------------------|
| a | = 9.49 (2) Å |
| b | = 8.48 (2) |
| c | = 7.71 (2) |
| α | = 107.8 (2)° |
| β | = 76.6 (2) |
| γ | = 100.1 (2) |
| V | = 565 (3) Å ³ |
| Z | = 1 |
| D_{exp} | = 1.80 g.cm ⁻³ |
| D_{calc} | = 1.82 |

Table 2. *Fractional coordinates, anisotropic thermal parameters and their standard deviations*

$$f_T = \exp \left[-\frac{1}{2}(B_{11}h^2a^*2 + B_{22}k^2b^*2 + B_{33}l^2c^*2 + 2B_{12}hka^*b^* + 2B_{13}hla^*c^* + 2B_{23}klb^*c^*) \right]$$

| | x/a | y/b | z/c | B_{11} | B_{22} | B_{33} | B_{12} | B_{13} | B_{23} |
|-------|------------|-------------|-------------|----------|----------|----------|-----------|-----------|----------|
| Ru | 0.1797 (1) | 0.1510 (1) | -0.0563 (1) | 2.33 (4) | 2.30 (4) | 2.64 (4) | 0.22 (5) | -0.81 (4) | 0.40 (4) |
| Cl(1) | 0.3549 (4) | 0.2344 (6) | 0.1412 (5) | 2.9 (1) | 4.0 (2) | 3.1 (1) | -0.6 (2) | -1.6 (1) | 1.0 (1) |
| Cl(2) | 0.0425 (4) | -0.0100 (5) | 0.1763 (4) | 2.4 (1) | 2.8 (1) | 2.5 (1) | -0.3 (2) | -1.1 (1) | 0.6 (1) |
| C(1) | 0.056 (2) | 0.354 (2) | 0.142 (2) | 3.7 (7) | 2.6 (5) | 3.0 (5) | 0.3 (8) | -0.1 (6) | 0.5 (5) |
| C(2) | 0.160 (2) | 0.420 (2) | 0.008 (3) | 2.6 (6) | 2.5 (5) | 4.8 (7) | 1.0 (7) | 0.9 (7) | 1.3 (6) |
| C(3) | 0.145 (2) | 0.342 (2) | -0.181 (2) | 5.3 (9) | 1.7 (4) | 2.7 (5) | 1.0 (9) | -1.4 (6) | -0.2 (4) |
| C(4) | 0.255 (2) | 0.375 (3) | -0.353 (3) | 3.8 (7) | 4.8 (10) | 3.8 (7) | 0.5 (11) | -0.7 (7) | 1.7 (8) |
| C(5) | 0.403 (3) | 0.326 (3) | -0.367 (3) | 4.4 (9) | 4.8 (10) | 5.4 (9) | -0.4 (12) | -0.4 (9) | 3.3 (9) |
| C(6) | 0.384 (2) | 0.185 (2) | -0.271 (2) | 3.0 (6) | 3.4 (6) | 3.3 (6) | 1.2 (8) | -1.2 (6) | 1.3 (6) |
| C(7) | 0.314 (2) | 0.024 (2) | -0.329 (2) | 1.8 (4) | 3.3 (6) | 3.2 (5) | 1.4 (6) | -0.1 (5) | -0.3 (6) |
| C(8) | 0.294 (2) | -0.082 (2) | -0.205 (2) | 2.4 (5) | 3.3 (6) | 3.9 (6) | 1.5 (7) | -0.1 (6) | 1.4 (6) |
| C(9) | 0.283 (2) | 0.558 (2) | 0.066 (3) | 2.8 (6) | 2.8 (6) | 6.0 (9) | -0.2 (7) | 0.1 (8) | 0.6 (8) |
| C(10) | 0.246 (2) | -0.038 (3) | -0.507 (2) | 4.3 (8) | 4.7 (8) | 2.8 (5) | -0.1 (12) | -1.8 (6) | -0.3 (6) |

Table 1 (cont.)

No absent reflexions
Space group $P1$ or $\bar{P}1$

Small changes of the above parameters from those reported in the preliminary communication (Porri *et al.*, 1965) are due to more accurate measurements.

Structure determination

From Patterson projections along the a and the c axes we were able to locate the Ru and Cl atoms; their posi-

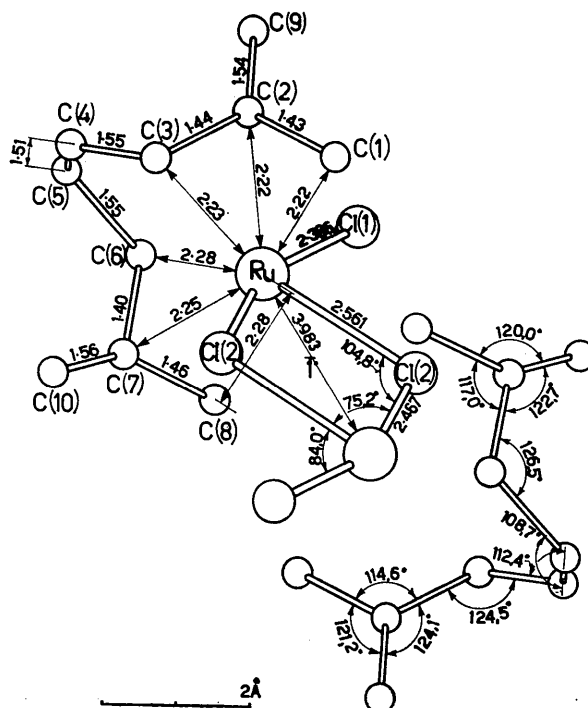


Fig. 1. Representation of the molecule seen along the crystallographic axis a . Bond distances and valence angles are reported. Standard deviations estimated for Ru-Ru, Ru-Cl Ru-C, C-C bonds respectively are 0.010, 0.007, 0.02, 0.03 Å; the first two σ 's are dominated by the e.s.d. of the unit cell parameters. For angles between carbon atoms $\sigma = 0.8^\circ$, for angles involving Ru and Cl atoms, $\sigma = 0.2^\circ$.

tions corresponded to a centrosymmetric arrangement, and hence the $P\bar{1}$ space group was accepted. Then by Fourier methods we were able to locate stepwise all carbon atoms on the same projections. After the refinement, carried out through differential Fourier syntheses, the $R = (\sum ||F_{\text{calc}}| - |F_{\text{obs}}|| / \sum |F_{\text{obs}}|)$ index was 0.09 and 0.13 for the non-zero $hk0$ and $0kl$ reflexions, respectively.

The atomic coordinates and thermal parameters obtained represented the starting point for the overall least-squares refinement, carried out with computer programs prepared by Immirzi (1967). The weighting scheme suggested by Cruickshank, Pilling, Bujosa, Lovell & Truter (1961) was adopted ($1/w = A + B|F_{\text{obs}}| + C|F_{\text{obs}}|^2$); a weight equal to half that corresponding to the weakest observed intensities was attributed to the non-observed reflexions. Atomic scattering factors were calculated with the analytic expression suggested by Vand, Eiland & Pepinsky (1957), adopting for the parameters the values suggested by Moore (1963).

Three cycles were run in the block-diagonal approximation, assuming anisotropic thermal parameters for the Ru and Cl atoms only. Then four cycles were performed with full-matrix and anisotropic thermal parameters for all non-hydrogen atoms; the final shifts of the atomic parameters were negligibly small, all of them being well below the corresponding estimated standard deviations. The final value of the R disagreement factor was 0.082 for the 1008 observed reflexions, and 0.105 if the 342 non-observed reflexions were accounted for with a fictitious structure factor equal to 0.5 times the limit of detectability. The final atomic parameters are reported in Table 2 with their standard deviations. Table 3 gives the full list of observed and calculated squared structure factors. Fig. 2 shows a perspective representation of the atomic thermal ellipsoids.

Discussion of the structure

The mode of coordination around the metal atom is

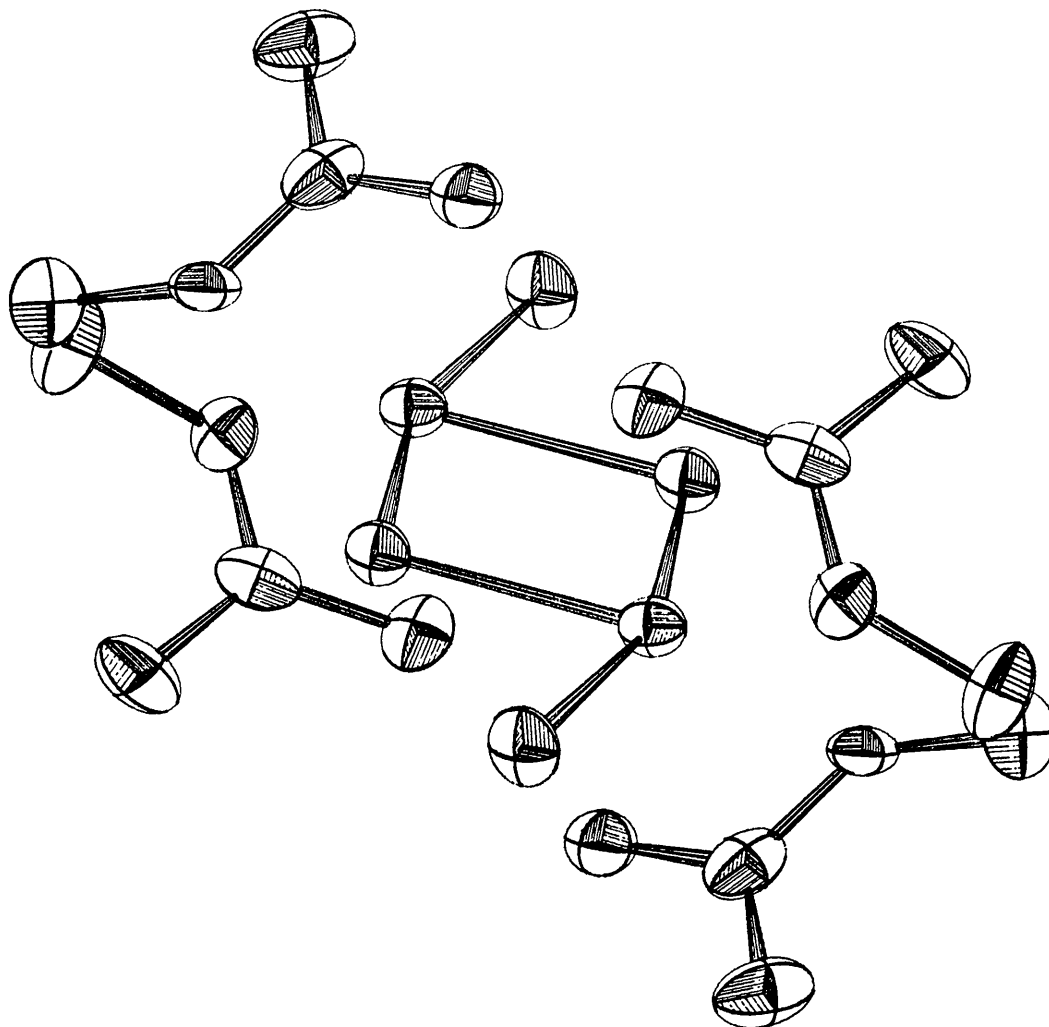


Fig. 2. Perspective representation of the atomic thermal ellipsoids, in the same projection as Fig. 1.

represented in Fig. 3. The resulting geometry is best described in terms of a slightly distorted trigonal bipyramid, provided we idealize each allyl group as a monodentate ligand, bonded to the metal through the centre of mass of the three allyl C atoms. In fact, the centres of the two allyl groups at the left- and right-hand side in Fig. 3 are only displaced by -0.07 and $+0.17$ Å with respect to the horizontal plane crossing

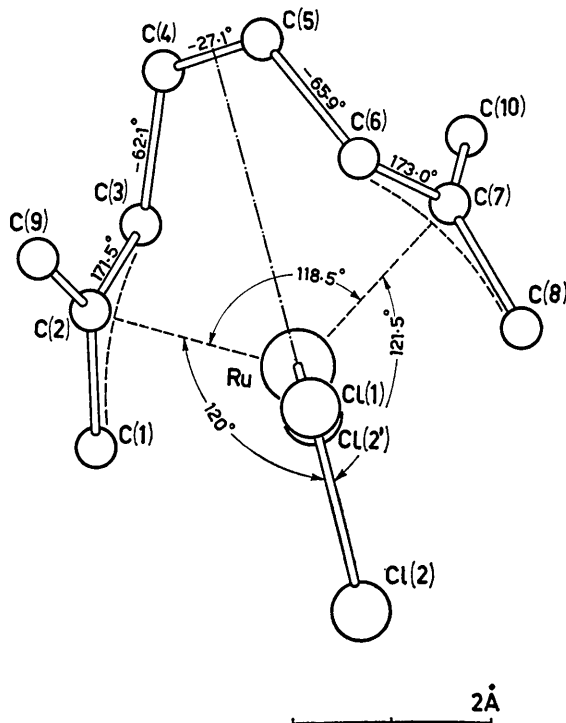


Fig. 3. The geometry of the coordination around the Ru atom is shown. Atoms Cl(2), C(2), C(7) constitute the equatorial vertices of the trigonal bipyramid, Cl(1) and Cl(2') (overlapping in the figure) its upper and lower vertices. The dihedral angles corresponding to the carbon skeleton of the ligand are indicated (*trans* conformation $\equiv 180^\circ$). The pseudo-twofold axis is indicated with a dash-and-dot line.

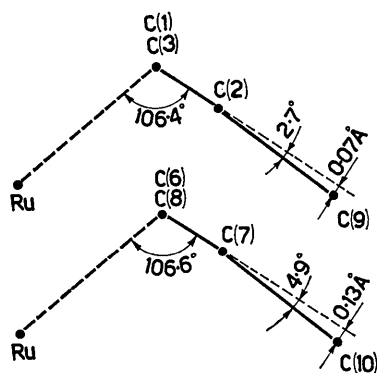


Fig. 4. The two allyl skeletons of the organic ligand seen along vectors C(1)-C(3) and C(6)-C(8).

the Ru atom; the corresponding displacement of the Cl(2) atom is -0.20 Å. The bond angles between the equatorial substituents are very close to 120° . The straight line joining the two apical chlorine atoms, Cl(1) and Cl(2'), is only 0.42 Å away from the metal atom, the distortion occurring toward the gap of the organic ligand, as in the case of $\text{RuCl}_2\text{C}_{12}\text{H}_{18}$ (Lydon & Truter, 1968). The molecular fragment shown in Fig. 3 contains a pseudo-twofold axis connecting the centre of the C(4)-C(5) bond and the ruthenium atom: the greatest atomic displacement from the exact twofold symmetry does not exceed 0.2 Å. The dihedral angles defined by the equatorial plane of Fig. 3 and the planes of the allyl groups are 65.4 and 63.3° for C(1)-C(2)-C(3) and C(6)-C(7)-C(8), respectively; the two values are close to the average value (63.2°) reported for $\text{RuCl}_2\text{C}_{12}\text{H}_{18}$ (Lydon & Truter, 1968).

As shown in Fig. 4, both methyl groups deviate from the allyl plane toward the half space containing the metal atom, in accordance with that reported for other metal complexes containing 2-methylallyl groups (Uttech & Dietrich, 1965; Mason & Russel, 1966; Mason & Wheeler, 1968). However, in contrast to that reported for the other cases, the angle between the C-CH₃ bond and the allyl plane is well below 20° . The two methylenic C atoms connected to the allyl groups are also slightly twisted out of the allyl planes, as is indicated by the C(1)-C(2)-C(3)-C(4) and C(5)-C(6)-C(7)-C(8) dihedral angles which deviate appreciably from 180° (see Fig. 3). The dihedral angle between the planes Ru, C(1), C(3) [or Ru, C(6), C(8)] and C(1), C(2), C(3) [or C(6), C(7), C(8)] is 106.4° (106.6°): these values are close, within a few degrees, to those reported for most metal-allyl π -complexes. The C-C bond lengths of the allyl groups do not show significant differences, nor do the Ru-C distances within each allyl group (see Fig. 1). One allyl group is significantly closer to Ru (2.22 Å being the average distance, with $\sigma = 0.012$ Å) than the other group (average distance = 2.27 Å, with the same σ). No explanation of this difference is possible on grounds of chemical bonding; on the other hand, the structure is centrosymmetric, and therefore no anomalous dispersion effect can account for the difference. The best explanation of this fact seems due to the different crystalline environment of the two allyl groups: as expected, the group which is closer to Ru is affected by more severe contacts with an adjacent molecule, consisting of three intermolecular C...C distances shorter than 3.6 Å (3.37 , 3.58 and 3.58 Å), while the other group is involved only in one distance in the same range ($d_{\text{C}\cdots\text{Cl}} = 3.53$ Å). On the whole, the Ru-C distances to coordinated carbon atoms are distributed in the same range as observed for $\text{RuCl}_2\text{C}_{12}\text{H}_{18}$.

As expected, the shortest Ru-Cl bond length (2.39 Å) corresponds to the simply coordinated Cl atom; this distance is remarkably lower than those reported for $\text{RuCl}_2\text{C}_{12}\text{H}_{18}$ (2.44 and 2.48 Å). The two Ru-Cl bridge bonds have different lengths (2.47 and 2.56 Å), the difference being well above the expected random fluctuation.

tuation ($\sigma_{\text{Ru-Cl}}=0.007 \text{ \AA}$). We believe that the greater bond length corresponding to the equatorially substituted chlorine atom is to be attributed to a repulsive effect on the Ru-Cl bonding electrons induced by the allylic π -electrons donated to the metal (see Fig. 3). Owing to the geometry of the coordination, the above effect is analogous to the well-known *trans* effect. Contrary to the case of most dimeric complexes containing chlorine bridges, the Ru_2Cl_2 planar four-membered ring deviates appreciably from the perfect square conformation: in addition to the above discussed difference in the Ru-Cl bond lengths, the Cl-Ru-Cl and Ru-Cl-Ru angles are quite different from 90° (75.2° and 104.8° respectively, see Fig. 1).

We express our gratitude to Drs Porri and Vitulli for providing us with the crystalline samples, and for very useful discussions.

Acta Cryst. (1971). B27, 1658

The Crystal Structure of 4-Ethylpyridinium Tetrabromoferrate(III)*

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(Received 1 July 1970)

Single-crystals of 4-ethylpyridinium tetrabromoferrate(III) have been prepared and their crystal structure has been determined by three-dimensional X-ray analysis. The compound crystallizes in the monoclinic space group $P2_1/c$ with $a=7.7068$ (8), $b=14.1673$ (11), $c=13.0414$ (16) Å, $\beta=84.19$ (1)°, and $Z=4$. Intensity data were measured using a θ - 2θ step scan technique on an automated diffractometer with Mo $K\alpha$ radiation. The structure was solved by superposition and Fourier methods and refined by full-matrix least-squares techniques to a final residual R of 0.061 for the 762 observed reflections used in the analysis. Ring atoms of the cation were refined as a rigid group. The FeBr_4^- tetrahedra are nearly aligned with their threefold axes almost coincident in the c direction. Average Fe-Br distance is 2.326 Å (this average becomes 2.347 Å when the interatomic distances are corrected for thermal motion assuming a riding model). The structure is in part stabilized by the presence of a weak N-H---Br hydrogen bond.

Introduction

The crystal-structure study of 4-ethylpyridinium tetrabromoferrate(III) was undertaken after the compound accidentally appeared as a byproduct during a series of investigations of bromo-coordinated antimony compounds in this laboratory. The true composition was determined from the crystal-structure analysis and also by an electron-microprobe analysis of the crystal used in the data collection. We felt it would be of interest to complete the crystal-structure analysis of the tetrabromoferrate(III) compound because of the general instability of iron(III)bromides (Sidgwick, 1950) and the

lack of crystal-structure data regarding the FeBr_4^- group.

Experimental

Crystal data

4-Ethylpyridinium tetrabromoferrate(III),
 $\text{C}_7\text{H}_9\text{NHFeBr}_4$,
 $M=483.65$, monoclinic $P2_1/c$, $F(000)=900e$, $Z=4$,
 $a=7.7068$ (8), $b=14.1673$ (11), $c=13.0414$ (16) Å,
 $\beta=84.19(1)^\circ$,
 $V=1416.6 \text{ \AA}^3$, $D_c=2.27 \text{ g.cm}^{-3}$,
 Mo $K\alpha$ ($\lambda=0.7107 \text{ \AA}$), $\mu=135.8 \text{ cm}^{-1}$.

Single crystals of $\text{C}_7\text{H}_9\text{NHFeBr}_4$ were obtained accidentally from contamination introduced into a vessel where the corresponding antimony bromide salt was being prepared. Of the two crystalline forms present, crystals of what later was determined to be the tetra-

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* Work was performed at the Ames Laboratory of the U.S. Atomic Energy Commission. Contribution No. 2777.

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